Experimental Report

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Abstract: Welding process induced in the heat affected zone (HAZ) and in the molten zone various complex phenomena: cyclic loading, miscrostructural evolutions, fluid flow within the weld pool, etc In many key industries such as transportation					

(aerospace, automotive...) and energy (nuclear) the prediction of residual stresses after welding is a challenge to know the service life of key components. To achieve this goal, the finite element (FE) method [1] is often used. Many commercial packages such as Abaqus, Code_Aster or Sysweld may be used to simulate the welding process. However when coupling between thermomechanical and microstructural effects is required as for age hardening materials, but the current FE models are not satisfactory. In fact, several phenomena that can occur in the HAZ are not classically considered: precipitates can nucleated, growth, coarse and dissolve [2][3][4][5]. These evolutions induce important differences on the mechanical behavior: during welding hardening or softening may occur, depending on the precipitation kinetics which is linked to the temperature history [3][4][5][6]. So a coupled approach (microstructural/mechanical) must be used to model the HAZ.

Experimental report: Internal stresses after welding for age hardening materials: advancedmodelisations vs neutrons characterizations

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Abstract

From an article that will be soon submitted. Here some experimental results are shown and the context is roughly presented.

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1. Neutrons diffraction at SALSA

A recent review presented in [1] gives an overview of the various methods to estimate the residual stresses. Here, our choice focuses on Neutrons Diffraction (ND) that can scan easily welded plates through the thickness while having an interesting resolution.

To get least biased possible residual strains, the ND were per-32 formed in the Laue Langevin Institut (ILL) in Grenoble on the $_{_{33}}$ 8 SALSA instrument [2]. This device uses a high neutrons flux $_{34}$ 9 and the measurements are not much affected by the background $_{35}$ 10 noise that is low on this instrument [3]. SALSA is a neutrons 11 diffraction instrument designed for the determination of strains 12 $\epsilon = (d_{hkl} - d_{hkl}^0)/d_{hkl}^0$ thanks to an accurate measurement of the ³⁶ 13 inter-reticular distances d_{hkl} which plays the role of gauge de-14 formation for a material under stresses [4]. 15

When a neutron flux is projected onto a material, each neu-³⁸ tron is diffused by the atoms of the target, these neutrons ³⁹ cause interference between them. For an orderly crystal we get ⁴⁰ builder interference for 2θ angles by respect to the Bragg's law: ⁴¹

$$n.\lambda = 2.d_{hkl}.sin\left(\theta\right) \tag{1}$$

To deduce from 2θ the average inter-reticular distance d_{hkl} of ⁴⁴ 16 a diffracting volume two kinds of methods are available: "poly-45 17 chromatic time of flight" and "monochromatic angular disper- 46 18 sive". The first is particularly suited for monitoring phase trans-19 formations, it consist to get a multitude of diffraction peaks 47 20 and the wavelength by gotten thanks the time of flight. The 48 21 second method, that is used on SALSA, is well suited to mea-49 22 sure the macroscopic stresses: only one diffraction peak 2θ is 50 23 recorded with a particular monochromatic wavelength [5] (here 24

*approx*1.64 Angstrom). In the case where a single diffraction peak is used, we choose for FCC alloys the (311) familly because they do not accumulate significantly inter-granular stresses during the sollicitation and therefore they have a representative behavior of the matrix [5, 4]. In addition, the 2θ peak is close to 90° that which provides a Gauge Volume (GV) that can be assimilated to a rectangular parallelepiped. From the diffraction peak, an average 2θ position on the GV is acquired and deformations or stresses can be deduced by comparison with an unstressed sample. Here, only first order stresses can be deduced: homogeneous on a grain aggregate.

2. Diffraction measurements

Generally the diffraction vector K (cf. fig. 2), which provides the measurement direction, focuses in three orthogonal directions to deduce the diagonal of the stress tensor. Here, the same approach is chosen, thus diffraction peaks are recorded in longitudinal, transverse and normal directions of our weld. Now, the choice of the GV depend on two main factors:

- the strain gradient which are expected in the structure. Here, the size of the MZ is about 0.5 mm close to weld foot. Then, a classical section 2x2 mm² for the GV should be avoided.
- the allocated beam time. Indeed, more the GV is low more the diffrated neutrons flux will be weak. However, to have a statistically good diffraction peak the count number on the detector has to be the highest possible.

For the measurements presented in this work, as well as the calibration/alignment step, the allocated time is 5 days for three differents plates. From the narrowness of the various zones in the weld, a $0.6x0.6 \text{ mm}^2$ section was chosen for the GV in

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normal and transverse directions (then the resolution is about 93 55 0.849 mm). This section being very low, the counting time to 56 obtain a reliable diffraction peak is large ($\approx 30 - 40$ min for 57 one direction). In addition, to capture strain gradients, recov-58 ery measurement was chosen (see fig. 2). So it was decided to 59 opt for a beam height of 20 mm in order to reduce the acqui-60 sition time. This dimension may seem important, nevertheless 61 in normal and transverse directions the height of the beam co-62 incides with the longitudinal direction of the weld joint (fig. 2). 63 Moreover, the neutrons diffraction has been applied in the cen-64 tral section of the plate and we have shown that in this area the 65 heat source has a stationary behavior. Thus we can expect a 95 66 quasi homogeneous mechanical behavior for a very reasonable 96 67 counting time (\approx 7 min per point). In the longitudinal direction, ₉₇ 68 the GV section is less restrictive but the height of the beam must 98 69 remain moderate. Therefore, it was reduced to 15 mm (which 99 70 is still important in this direction) and a 2x2x15 mm³ GV was₁₀₀ 71 chosen (fig. 2). 72 101



Figure 1: Example of resultats far from the weld for a plate welded with V = 118 0.45 m/min and V = 0.9 m/min with a raw peak reference.

The measurements gotten on the 2D SALSA detector were₁₂₁ 73 analyzed by using the bean line software: LAMP [6]. The₁₂₂ 74 diffracted intensities I are integrated to obtain an unidirectional₁₂₃ 75 diffraction peak $(I,2\theta)$ as show in fig. ??. Following a correc-124 76 tion of the background noise (a linear correction is sufficient in 77 the presence of a sharp peak. For others peaks a noise func-78 tion is deduce) a Gaussian fitting is performed and its center 79 gives the 2θ position. In a majority of results, the error bar pro-80 vided by the adjustment procedure (given by LAMP) is very 81 low. However, in the longitudinal direction, the first measure-82 ment shown lows diffracted signals. It has been therefore de-83 cided to carry out shorter steps in this direction but with three 84 different rotation angles (±1.5°) for a same diffraction direc-85 tion and then summed. Although a significant improvement 86 has been noted, the quality of the diffraction signal obtained re-87 mains lower quality compared to the transverse and normal di-88 rections. This effect cannot be attributed to a texture effect, then 89 it is probably due to a too low count time or an columnar grain 90 influence which provides a lower amount of diffracted grains. 91 To compensate these effects and get good quality results, a four-92

point moving average has been effected .

3. Residual elastic strains results

To get the residual elastic deformations $\epsilon = \Delta d_{hkl}/d_{hkl}^0$ in the material the variation of inter-reticular distance (or diffraction peak) must be performed. On SALSA, the wavelength λ can be assumed constant, then we have thanks to a Bragg's law derivate:

$$\epsilon = -\frac{\Delta 2\theta}{2.tan(\theta)} \tag{2}$$

where $\Delta 2\theta$ is the peak displacement (from $2\theta_0$) for the same alloy without stresses and with the same local composition. Of course, this is an idealized vision because it is very difficult to satisfy these requirements in a strict way. To have a good reference value $\Delta 2\theta_0$ and to deduce *Delta* 2θ , measurements were performed on a superposition of thin "combs" (thickness 1 mm and width 1.5 mm), which were cut from a section in the welded plate (see fig. 3). This structure allows to relax residual stresses to get a reference that best approximates a mechanical state without stresses. Then, measurements were performed in the MZ, HAZ and the BM of the combs to overcome the chemical influence. During these measurements, it was observed an expansion of the lattice from the BM to the HAZ (a dissolution of precipitates induces in the matrix a homogeneous presence of many solutes atoms) and a slight compression from the HAZ to the MZ (probably due to the smaller grain size which traps many atoms and intermetallic and therefore generates less intragranular expansion).

These references were therefore applied as a constant in the MZ and BM areas and a spline interpolation has been used in all peaks positions in the HAZ.

The first results in fig. 1 shows away from the weld (distance about 66 mm) that the residual deformations are not zero in the Base Metal. This effet can be caused by rolling process but also by a non-zero stress state in the reference combs (cf. fig. 3). The numerical simulations presented in the next section suggests that this residual strain state away from the weld must be zero. Thus, a weak shift on the diffraction peak (0.5% of the current value) was performed to have in average 0% far from the weld.



Figure 2: (a) Illustration of measured points and (b) representation of the several Gauge Volume for each orientations.

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Figure 3: Combs used for the reference and representation of the chemical effect on the peak position 2θ .

Now, the figs. 4, ?? and ?? show the results of residual elas-125 tic strains in three different directions for two plates: V =126 0.45 m/min and V = 0.9 m/min (the results for the plate 127 V = 0.72 m/min are not show because they are between the 128 two others). It is important to note that the error bars presented 129 in these figures represent only the error caused by the adjust-130 ment procedure and not the various possible sources of errors. 131 Indeed, it is just the fitting error provided by LAMP software. 132 So it does not take into account these uncertainties: (i) even-133 tual defects in alignment, (ii) microstructure "defects", (iii) er-134 ror caused by the instrumental resolution, and (iv) inaccuracies 135 during reference measurement (potentially a complex state of 136 residual stresses in the comb). 137

In view of the results on the transverse lines (measurements 138 in the thickness direction are not presented), one can first be sat-170 139 isfied with a relative good symmetry on both sides of the weld 140 center. Then, the strain amplitudes are close to zero in the nor-141 mal direction, except in the vicinity of the HAZ-BM transition.173 142 In the longitudinal direction, the strain is near zero away from¹⁷⁴ 143 the HAZ and it increases progressivly close to the MZ and then¹⁷⁵ 144 decreases in the center. Finally, for the transverse direction, we 145 have a weak compression state in the weld center which be-178 146 comes a tension state in the HAZ-BM border and then decline¹⁷⁹ 147 180 close to zero far away. 148

These results are close to the behavior observed during the laser welding process of thin plates gotten thanks to finite ele-183 ment simulations for a 6056 aluminum alloy [7].

152 4. Residual stresses estimation

Strain measurement have been carried out in three orthogo-153 nal directions in several areas of the weld. This choice allows 154 to obtain an estimation of residual stresses amplitude in the ma-155 terial. In fact, by using the Hooke's law it is possible to get 156 the diagonal of the stresses tensor σ_{ii} with the elastic Lame co-157 efficients (λ, μ) thanks to longitudinal, transverse and normal 158 residual strain experiments $\sigma_{ii} = \lambda \epsilon_{kk} \delta_{ii} + 2\mu \epsilon_{ii}$. In this case, 159 the shear components cannot be determined. 160

It should be noted that this step (the summation by Hooke's law) does not occur without source of mistakes: (i) the using of Lame coefficient with the (311) strains neglect the anisotropy of the FCC structure [5], (ii) the summation mixte unaccurate measurements (here longitudinal) with accurates ones and then all σ_{ii} components are impacted by the worse direction, (iii) furthermore we sum results with two different resolutions and dissimilar GV. This is why, in this work, only the residual elastic strains are compared in fig. 4 to the simulations.



Figure 4: Residual elastic deformation for a transverse weld line at 4 mm from the lower surface for the transverse direction. Confrontation between numerical and experimental results for a welding velocity of V = 0.45 m/min. Additionnal experimental results are shown with V = 0.9 m/min to quantify the velocity influence.

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