Proposal:	4-02-440			Council: 10/2014			
Title:	Mechanism of the antiferrom	hanism of the antiferromagnetic transition in the geometrically frustrated					
Research area: Physics							
This proposal is a new proposal							
Main proposer:	Sylvain PETIT						
Experimental te	eam: Elsa LHOTEL						
	Solène GUITTENY						
	Sylvain PETIT						
Local contacts:	Hannu MUTKA						
	Jacques OLLIVIER						
Samples: Er2Ti	207						
Instrument		Requested days	Allocated days	From	То		
IN5		5	6	15/07/2015	21/07/2015		
Abstract:							
We would like to stu pyrochlore Er2Ti2O	dy the evolution of the spin 7. The aim of the experiment	gap as a function of the time time the time time the time time time time time time time tim	f the temperature a mechanism respo	and of magnetic finds	eld in the geometrically fru pilization of the long range	ustrated order	

pyrochlore Er2Ti2O7. The aim of the experiment is to determine the mechanism responsible for the stabilization of the long range order in this material. Indeed, recent reports claim that it might be "order by disorder", but the CEF small anisotropies might also be at play. We hope to be able to get a set of results allowing to discriminate between these two scenarii.

Experimental report on Exp 4-02-440

Geometrical frustration and the associated strong degeneracy of the ground state manifold very often prevent the stabilization of standard magnetic phases. At this point, quantum or thermal fluctuations enter into play to select and stabilize a particular configuration (or a subset of configurations), a phenomenon called "order by disorder" [2]. It has been recently argued that the pyrochlore compound $Er_2Ti_2O_7$ is a model system to study these subtle effects [3,4,5] (in $Er_2Ti_2O_7$, the rare-earth magnetic moments are localized at the vertices of corner-sharing tetrahedra. They present a strong XY-like anisotropy, the XY planes being perpendicular to the local <111> ternary axes).

Actually, $Er_2Ti_2O_7$ undergoes a transition towards an antiferromagnetic k=0 Néel phase below T_N =1.2K [5,6,7,8]. This magnetic ordering was described in detail by neutron polarimetry (ψ 2 vector in Γ 5 representation), as a non-collinear antiferromagnetic structure, with magnetic moments, perpendicular to the local <111> axes, and aligned along <211> [9].

Several theoretical works have proposed that this peculiar ordering can be explained by a quantum order-by-disorder mechanism [10,11]. This model is based on a pseudo spin ½ Hamiltonian, spanning the ground crystal field doublet. It is characterized by an extensive degeneracy which is relieved by quantum zero point fluctuations contribution out of the ψ 2 configuration. In a recent publication, we have proposed another mechanism called "order by virtual crystal field excitation" based on a mean field model written in terms of the actual Er^{3+} magnetic moment and taking into account the full CEF scheme (Ref [12,13]). Here, the ψ 2 configuration is selected by a small anisotropy due to the exchange induced mixing between the CEF states. We have shown that the spin dynamics calculated in both models compare quite well with the neutron data recently collected at IN5 (see [13] as well as figure 1). Further, both models predict the opening of a spin gap at zone centers. The data of Ref [14] along with previous IN5 experiment confirm this prediction (see [13] and Figure 2), even if the experimental gap is larger, 43 µeV instead of 21 and 15 in the two above models.

Unfortunately, the two models remain quite difficult to distinguish. As a further possible test, we proposed in the present experiment to apply a magnetic field, and see how the spin gap evolves. With 8.5 Å the energy resolution was sufficient to observe the spin gap at the Q=(111) (see Fig. 4) between 0 and 1.5 T. The spin gap as a function of field measured at Q=(111) is then shown in Figure 3–left.





The prediction of our model "order by virtual crystal field excitation" corresponds to the open green points in Figure 3-right. The interpretation of the data is still in progress.

References

[1] C. Lacroix, Introduction to Frustrated Magnetism, edited by C. Lacroix, P. Mendels, and F. Mila (Springer-Verlag, Berlin, 2011).

- [2] J. Villain, R. Bidaux, J.-P. Carton, R. Conte, J. Phys 41, 1263 (1980).
- (3] J. S. Gardner, M. J. P. Gingras, J. E. Greedan, Rev. Mod. Phys. 82, 53 (2010).
- [4] S. T. Bramwell and M. J. P. Gingras, Science 294, 1495 (2000).
- [5] J. D. M. Champion et al Phys. Rev. B 68, 020401 (R), (2003)
- [6] W. J. Blote, R.F. Wielinga and W. J. Huiskamp, Physica 43, 549 (1969).

[7] M. J. Harris, S. T. Bramwell, T. Zeiske, D. F. McMorrow, and P. J. C. King, J. Magn. Magn. Mater. 177, 757 (1998).

[8] R. Siddharthan, B. S. Shastry, A. P. Ramirez, A. Hayashi, R. J. Cava, and S. Rosenkranz, Phys. Rev. Lett. **83**, 1854 (1999). [9] A. Poole, A. Wills, E. Lelievre-Berna JPCM **19**, 452201, (2007).

[10] M. E. Zhitomirsky, M. V. Gvozdikova, P. C. W. Holdsworth, and R. Moessner, Phys. Rev. Lett. 109, 077204 (2012).

- [11] L. Savary, K. A. Ross, B. D. Gaulin, J. P. C. Ruff and L. Balents, Phys. Rev. Lett. 109, 167201 (2012).
- [12] P. A. McClarty, S. H. Curnoe and M. J. P. Gingras, Journal of Physics: Conference Series 145, 012032 (2009).
- [13] S. Petit, J. Robert et al, Phys. Rev. B **90**, 060410(R) (2014)
- [14] K. A. Ross, Y. Qiu, J. R. D. Copley, H. A. Dabkowska, and B. D. Gaulin, Phys. Rev. Lett. 112, 057201 (2014).

Figure 3 : raw data taken at 50 mK along (11L) for various fields.

