

Experimental report

06/09/2018

Proposal: 5-32-853

Council: 4/2017

Title: Neutron diffraction study of Sr₂FeCoO₆, SrLaFeCoO₆ and La₂FeCoO₆

Research area: Materials

This proposal is a new proposal

Main proposer: Tapan CHATTERJI

Experimental team: Tapan CHATTERJI

Local contacts: Lucile MANGIN-THRO

Samples: La₂FeCoO₆
SrLaFeCoO₆
Sr₂FeCoO₆

Instrument	Requested days	Allocated days	From	To
D7	7	3	15/06/2018	18/06/2018

Abstract:

Double perovskites are a fascinating material class with a lot of functional properties. Multiferroicity, magnetocapacitance, magnetoresistance, different magnetic ground states including spin-glass and three dimensional quantum magnetism have been either experimentally observed or theoretically proposed for these materials. We propose to do polarized neutron diffraction study on the double perovskites Sr₂FeCoO₆, SrLaFeCoO₆ and La₂FeCoO₆ as a function of temperature in the range 1.5 – 300 K in fine temperature steps. The XYZ polarization analysis available on the polarized diffractometer D7 will enable us to separate magnetic and nuclear scattering absolutely essential in order to interpret the data properly. We request 7 days of neutron beam time on D7 for this.

Polarized neutron diffraction study of $\text{Sr}_2\text{FeCoO}_6$, SrLaFeCoO_6 and $\text{La}_2\text{FeCoO}_6$

Introduction:

Double perovskites are a fascinating material class with a lot of functional properties. Multiferroicity, magnetocapacitance, magnetoresistance, different magnetic ground states including spin-glass and three dimensional quantum magnetism have been either experimentally observed or theoretically proposed for these materials. One of the famous double perovskite discovered is the $\text{Sr}_2\text{FeMoO}_6$, where high magnetoresistance (MR) was reported [1]. A strong correlation between the observed MR and the inherent crystallographic *antisite* disorder was proposed in that compound. The *antisite* disorder – where the transition metal atoms occupy the crystallographic positions interchangeably – is found to be an ubiquitous presence in most of the double perovskite compounds and plays an important role in the magnetism of these materials. Theoretically proposed (using the Goodenough-Kanamori rules) to behave as long-range ordered magnets, the double perovskites often display strong signs of disordered magnetism or short-range magnetic order primarily due to the presence of *antisite* disorder. Following the work on $\text{Sr}_2\text{FeMoO}_6$, we extended our research to related compounds $\text{Sr}_2\text{FeCoO}_6$, SrLaFeCoO_6 and $\text{La}_2\text{FeCoO}_6$ [2-4].

In order to achieve a better understanding of the disorder effect and spin states of Fe and Co, a suitable non magnetic doping at the *A* site can be done. *A*-site doping in double perovskites have shown that *antisite* defect have a role to play [5]. Since Fe and Co have a similar x-ray scattering factor, it is necessary to do neutron diffraction analysis to determine the amount of *antisite* disorder. Also, it has been known that the degree of disorder can be understood better by analyzing the variation of anisotropic thermal parameters with respect to temperature [6]. Substituting La at Sr-site in $\text{Sr}_2\text{FeCoO}_6$ modifies the magnetic and transport properties. An interesting change is that, while $\text{Sr}_2\text{FeCoO}_6$ showed a clear spin-glass nature in dynamic susceptibility studies such a behavior is not observed in SrLaFeCoO_6 . Though there is irreversibility seen in dc magnetization, no frequency dependence of the transition is observed in ac susceptibility. Another remarkable fact is the hysteresis observed in the real part of ac susceptibility. Though the MR gets reduced by 10%, a hysteretic nature in MR is observed at 5 K, which seems to vanish at 50 K. The magnetization is found to increase with La substitution, which is quite the opposite that has been found in $\text{Sr}_2\text{FeMoO}_6$.

On the other extreme, when Sr is completely replaced by La it shows no magnetoresistance or spin-glass behaviour. Interestingly, at room temperature, $\text{La}_2\text{FeCoO}_6$ crystallizes in rhombohedral *R-3c* space group. As the temperature decreases, it is seen to develop a mixed-phase consisting of *R-3c* and orthorhombic *Pnma*. The presence of this structural mixed-phase reflects in magnetic measurement as a thermal hysteresis. Further, a broad magnetic phase transition is seen to occur at temperatures very close to 300 K.

Experimental results:

We have done the polarized neutron diffraction experiment on D7 on the double perovskites. In the limited allocated neutron beam time of only three days allocated by the subcommittee we could measure only two compounds; out of these $\text{La}_2\text{FeCoO}_6$ which was suspected to order magnetically, was measured rather well for two and half days and the other magnetic glass compound SrLaFeCoO_6 could be measured only partially for half a day. The interesting canonical spin-glass compound $\text{Sr}_2\text{FeCoO}_6$ could not be measured at all. We have done XYZ polarization analysis of the neutron diffraction intensities and separated them into nuclear coherent, nuclear spin incoherent, spin flip magnetic and non-spin flip magnetic parts. Here we have considered only the average magnetic intensities in the following results.

Fig. 1 (a) shows 011 magnetic peaks at several temperatures of $\text{La}_2\text{FeCoO}_6$ that appear at about $Q = 1.4 \text{ \AA}^{-1}$ below about $T_N = 220 \text{ K}$. We have fitted these peaks with a Gaussian function. Fig 1(b) shows the temperature variation of intensity of 011 reflection and its power-law fit. The fit gives $T_N = 250 \pm 17 \text{ K}$ and exponent $\beta = 0.44 \pm 0.07$. The appearance of 011 magnetic peak below $T_N = 250 \text{ K}$ suggests G-type magnetic ordering at lower temperatures.

For SrLaFeCoO_6 we have observed no magnetic Bragg peaks. But we have observed diffuse magnetic scattering close to $Q = 1.4 \text{ \AA}^{-1}$ and $Q = 1.6 \text{ \AA}^{-1}$. Fig. 2 (a) shows the magnetic diffuse scattering observed in the Q range $1 - 2 \text{ \AA}^{-1}$. We have fitted two Lorentzian functions for the two peaks. The peaks at $Q = 1.4$ and 1.6 \AA^{-1} represents spin fluctuations. We are still analyzing the data to understand its implications. We have also observed temperature

dependence of the background close to $Q = 0$ shown in Fig. 2 (b). The higher background at low temperatures may signify ferromagnetic correlation at $Q = 0$. The data for SrLaFeCoO_6 , are preliminary only. To do any definite conclusions about the short-range correlation and its temperature dependence we need to measure this sample with much higher counting statistics.

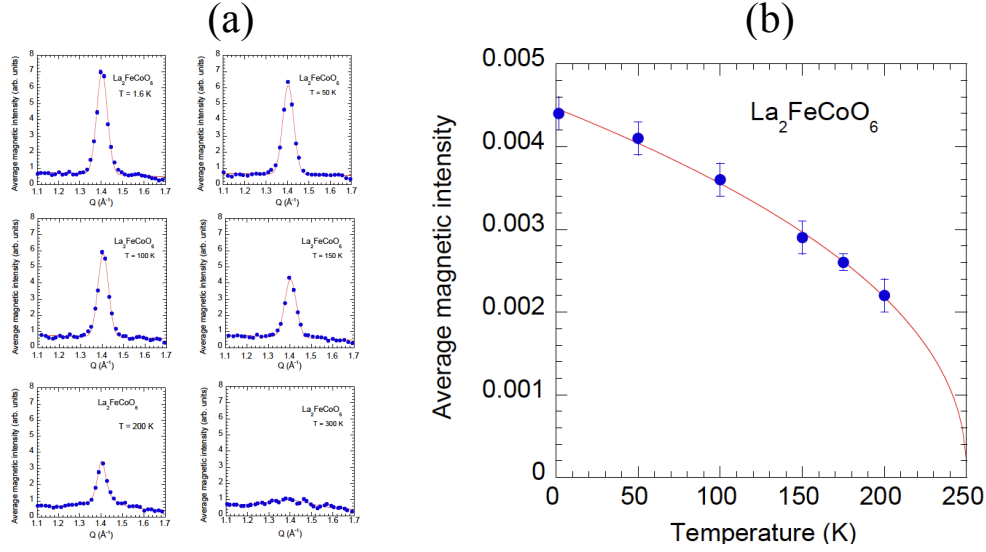


Fig. 1 (a) The 011 magnetic peaks at several temperatures of $\text{La}_2\text{FeCoO}_6$ that appear at about $Q = 1.4 \text{ \AA}^{-1}$ below about $T_N = 220 \text{ K}$. We have fitted these peaks with a Gaussian function. (b) Temperature variation of intensity of 011 reflection and its power-law fit. The fit gives $T_N = 250 \pm 17 \text{ K}$ and exponent $\beta = 0.44 \pm 0.07$.

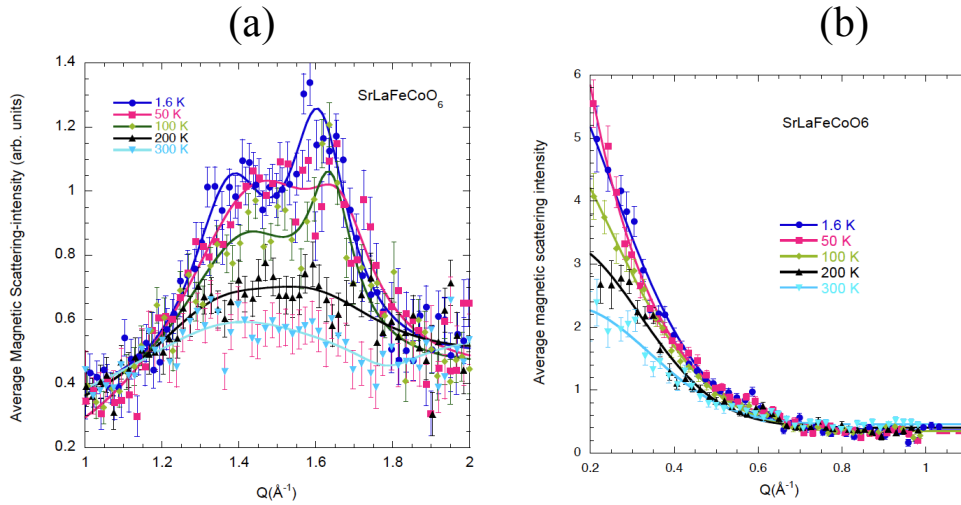


Fig. 2 – (a) Magnetic diffuse scattering from SrLaFeCoO_6 . We have apparently two magnetic broad peaks which we fitted with two Lorentzians. The widths of these peaks are much larger than the resolution width of the instrument. These broad peaks indicate short range magnetic order. (b) Temperature variation of the background of from SrLaFeCoO_6 close to $Q = 0$. The higher background at low temperatures may signify ferromagnetic correlation at $Q = 0$.

References:

- [1] K.-I. Kobayashi et al., *Nature* **395**, 677 (1998); D.D. Sarma et al., *Phys. Rev. Lett.* **85**, 2549 (2000); D.D. Sarma et al., *Solid State Comm.* **114**, 465 (2000); K.-I. Kobayashhi et al., *Phys. Rev. B* **59**, 11159 (1999).
- [2] Pradheesh R et al., *J. Appl. Phys.* **111**, 053905 (2012).
- [3] Pradheesh R et al., *Appl. Phys. Lett.*, **101**, 142401, 2012
- [4] Pradheesh R et al., *Eur. Phys. J. B*, **85**, 260, 2012.
- [5] J.Navarro, J.Nogués, J.Munoz, and J.Fontcuberta, *Phys. Rev. B* **67**, 174416 (2003).
- [6] P. G. Radaelli et al., *Phys. Rev. B* **54**, 8992 (1996)