

Experimental report

08/02/2024

Proposal: 5-41-1211

Council: 4/2023

Title: Magnetic Structure of the spin-1/2 antiferromagnetic XXZ chain PbCo₂V₂O₈ in a transverse (1,1,0) magnetic field

Research area: Physics

This proposal is a new proposal

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Samples: PbCo₂V₂O₈

Instrument	Requested days	Allocated days	From	To
D23	6	4	29/10/2023	02/11/2023
ORIENTEXPRESS	1	1	06/11/2023	07/11/2023

Abstract:

The one-dimensional (1D) spin-1/2 antiferromagnetic with Heisenberg-Ising (XXZ) anisotropy, is a fascinating model system to explore fundamental physics such as quantum phase transitions and exotic ground states and excitations. Although no long-range magnetic order is expected in the pure 1D case, real materials which have interchain coupling develop long-range magnetic orders that reflect the nature of these ground states. The ACo₂V₂O₈ compounds (A=Sr, Ba) are great physical realizations of the Heisenberg-Ising spin-1/2 chains. They show exotic physics in zero field as well as [1,0,0] and [0,0,1] applied fields with complex magnetic phases and excitations that can be used to test physics concepts. We have discovered a previously unreported phase for the [1,1,0] transverse field direction that is accessible in only the newly synthesized PbCo₂V₂O₈ member of this family due to its lower energy scale. The aim of this experiment is to determine the magnetic structure of this field-induced phase using single crystal neutron diffraction.

EXPERIMENTAL REPORT

ON EXPERIMENT PERFORMED AT JCNS INSTRUMENT AT ILL OR SNS

Proposal number:	5-41-1211		
Experiment title:	Magnetic Structure of the spin-1/2 antiferromagnetic XXZ chain $\text{PbCo}_2\text{V}_2\text{O}_8$ in a longitudinal (0,0,1) magnetic field		
Instrument	D23		
Dates of experiment:	29/10/23 – 02/11/23	Date of report:	08/02/24
Experimental team:			
Names	Addresses		
	Cintli Aguilar Maldonado, Konrad Puzniak, Ralf Feyerherm, Nazmul Islam, Bella Lake		
Local Contact:	Eric Ressouche, Karin Schmalzl		

Abstract

The one-dimensional (1D) spin-1/2 antiferromagnetic with Heisenberg-Ising (XXZ) anisotropy, is a fascinating model system to explore fundamental physics such as quantum phase transitions and exotic ground states and excitations. Although no long-range magnetic order is expected in the pure 1D case, real materials which have interchain coupling develop long-range magnetic orders that reflect the nature of these ground states. The $\text{ACo}_2\text{V}_2\text{O}_8$ compounds ($A=\text{Sr}, \text{Ba}$) are great physical realizations of the Heisenberg-Ising spin-1/2 chains. They show exotic physics in zero field as well as $[1,0,0]$ and $[0,0,1]$ applied fields with complex magnetic phases and excitations that can be used to test physics concepts.

Introduction

One dimensional (1D) spin-1/2 antiferromagnets with Heisenberg-Ising (XXZ) anisotropy have been a topic of recent study, in particular using the $\text{ACo}_2\text{V}_2\text{O}_8$ family ($A=\text{Sr}, \text{Ba}$) which represent model materials to explore fundamental physics. The characteristic 4-fold screw chains of magnetic ions (Co^{2+}) along its tetragonal c-axis with effective spin- $\frac{1}{2}$ exhibit a strong antiferromagnetic intrachain coupling, while the weak interchain coupling gives rise to long-range, antiferromagnetic Neel order at low temperatures. In zero magnetic field these compounds have a spinon continuum above the Neel temperature $T_N \sim 5 \text{ K}$ which is used to demonstrate spinon confinement on cooling below T_N where the continuum is replaced by sharp bound-spinon modes $[1,2]$.

Exotic physics arises as a magnetic field is applied parallel to the easy axis (c-axis) namely a longitudinal field. The Neel antiferromagnetic order (NAF) is suppressed to much lower values ($T_N < 1$ K) and above a critical field $\mu_0 H_{c1}$ the system undergoes a transition to a longitudinal spin density (LSDW) and then another transition at $\mu_0 H_{c2}$ to a transverse canted antiferromagnet (TAF) [3,4]. The phase diagram for the longitudinal magnetic field versus temperature was found by several techniques including neutron diffraction, NMR and specific heat for $\text{BaCo}_2\text{V}_2\text{O}_8$, and $\text{SrCo}_2\text{V}_2\text{O}_8$ [5,6]. Neutron diffraction revealed a commensurate to incommensurate transition at $\mu_0 H_{c1}$ above which the magnetic reflections can be indexed by propagation vector $k_{IC} = (0, 0, 1 \pm \delta l)$, while the magnetic order above $\mu_0 H_{c2}$ is also incommensurate along the c-axis. Within these field-induced phases spinon pair excitations (psinons) and bound states of magnons (Bethe Strings) have been discovered and explored using inelastic neutron scattering and terahertz spectroscopy [6,7].

We have synthesized a new family member namely $\text{PbCo}_2\text{V}_2\text{O}_8$, and powder measurements showed it to have very similar magnetic properties to $\text{SrCo}_2\text{V}_2\text{O}_8$ and $\text{BaCo}_2\text{V}_2\text{O}_8$. The magnetic structure in zero field is the same for all three family members [8] and the H - T phase diagram for longitudinal and transverse field parallel to $[1,0,0]$ are similar except that $\text{PbCo}_2\text{V}_2\text{O}_8$ has a lower energy scale. The H - T phase diagram of $\text{PbCo}_2\text{V}_2\text{O}_8$ as a function of temperature and longitudinal magnetic field (the phases have to be confirmed) is presented in fig. 1 below.

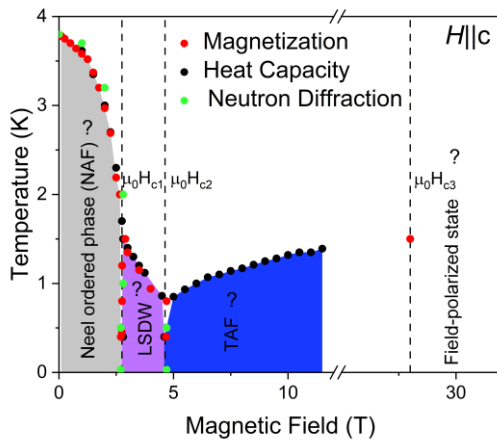


Figure 1. Phase diagram of $\text{PbCo}_2\text{V}_2\text{O}_8$ as a function of longitudinal magnetic field $H||c$ and temperature. The different colored regions indicate the different magnetic phases. The gray phase is identified as a Néel antiferromagnetic phase. The purple phase might have longitudinal spin density wave order, whereas the blue phase might be a transverse antiferromagnetic phase, although these field-induced phases have yet to be confirmed [8].

The phase diagram reveals a suppression of the Neel ordering temperature and transitions to two field-induced phase at $\mu_0 H_{c1} = 2.7$ T and $\mu_0 H_{c2} = 4.6$ T with increasing field. Although the nature of these phases is as yet unknown we anticipate that they are similar to those of $\text{SrCo}_2\text{V}_2\text{O}_8$ and $\text{BaCo}_2\text{V}_2\text{O}_8$.

Sample Details and Instrumental Configurations

The neutron diffraction measurements were collected, below T_N at 30 mK and the single crystal neutron diffraction was performed on a small cylindrical-shaped crystal (weight 1 g, diameter 5 mm, and length 6 mm) at the D23 neutron diffractometer. The single crystal was previously aligned on the Orient Express neutron Laue diffractometer. The sample was then installed on D23 in a vertical field orange cryomagnet (equipped with a dilution insert) with the c-axis aligned along the magnetic field. A maximum longitudinal magnetic field of 13 T ($H||c$) could be reached and the incident wavelength of the neutron beam is $\lambda = 1.28$ Å.

Objectives

The main goal of this experiment was to determine the magnetic structure of $\text{PbCo}_2\text{V}_2\text{O}_8$ in a longitudinal $[0,0,1]$ magnetic field.

Experimental Results

Powder and single crystalline samples of $\text{PbCo}_2\text{V}_2\text{O}_8$ have been studied in our research group where it has been proved to be isostructural to $\text{SrCo}_2\text{V}_2\text{O}_8$ with space group $I41cd$ (No. 110). $\text{PbCo}_2\text{V}_2\text{O}_8$ develops long-range Neel antiferromagnetic order below $T_N = 3.8$ K, its magnetic structure at 0 T is the same as the other members of the $\text{ACo}_2\text{V}_2\text{O}_8$ family with a propagation vector $q = (0, 0, 1)$. The spins point mostly parallel to the tetragonal

c-axis and align antiferromagnetically along the screw chains following the Γ_5 irreducible representation [8]. The aim of this experiment was to determine the magnetic structure of field-induced phases using single crystal neutron diffraction. For this purpose, a single crystal of $\text{PbCo}_2\text{V}_2\text{O}_8$ was prepared and cut to appropriate dimensions for the experiment. The sample was aligned with the $[0,0,1]$ direction vertical and parallel to the magnetic field direction and glued with X60 glue into a copper sample holder with the pertinent size for the 15 T vertical magnet with ^3He insert experimental setup. A copper monochromator was required to provide a neutron wavelength of $1.28 \text{ \AA}$ and the ^3He insert was also required to achieve temperatures down to 30 mK.

As reported for $\text{SrCo}_2\text{V}_2\text{O}_8$ a commensurate to incommensurate transition occurs at $\mu_0 H_{c1} = 3.9 \text{ T}$ where the LSDW phase begins and given the similar behavior in all members of the family, $\text{PbCo}_2\text{V}_2\text{O}_8$ is expected to have the same behavior at $\mu_0 H_{c1} = 2.7 \text{ T}$. A field dependence of the magnetic reflection $(2, 3, 0)$, up to 6 T, is shown in fig 2 [9]. The peak broadens, however at higher fields, above 2.7 T, we could not observe the two satellite reflections: $(2, 3, \delta l)$ and $(2, 3, -\delta l)$ expected if there is incommensurate order [4,10]. One of the possible explanations for this is due to the experimental setup, which did not allow us to have the best resolution along c direction (since the 15 T vertical magnet has a very limited vertical access of -2 to 2 degrees and no collimator was used) to resolve the splitting. It should be emphasized that the focus of the previous experiment was to measure the $[1,1,0]$ field direction [11], thus the setup was not optimized for a measurement along $[0,0,1]$ field direction. A followup experiment using a vertical collimator along with a vertical magnetic that has wide vertical angular access, or alternatively a horizontal magnet would clarify the situation.

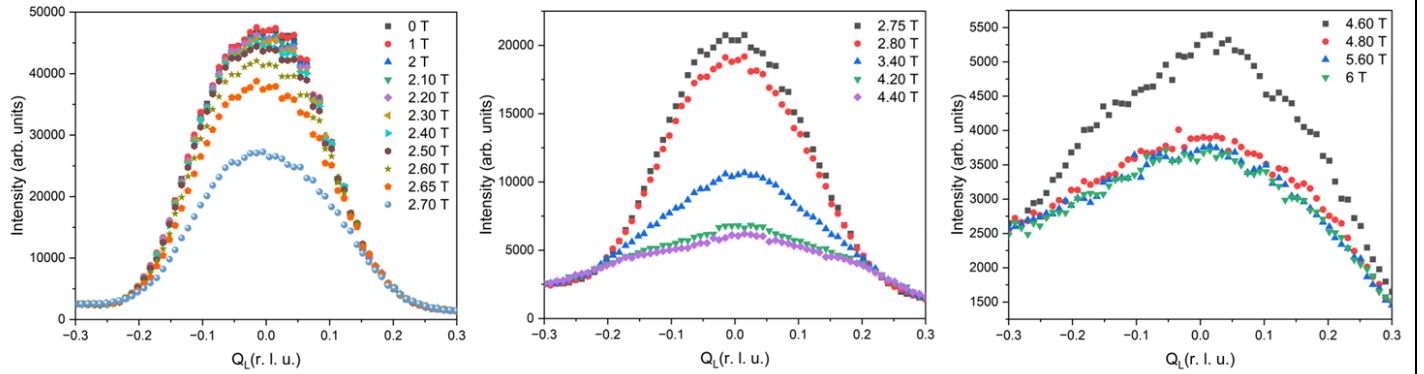


Figure 2. Field dependence of magnetic reflection $(2,3,0)$ up to 6 T [9].

References

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