

<b>Proposal:</b>	5-41-698	<b>Council:</b>	4/2012
<b>Title:</b>	Magnetic structure of SmFeO <sub>3</sub>		
<b>This proposal is a new proposal</b>			
<b>Research Area:</b>	Physics		

<b>Main proposer:</b> KOMAREK ALEXANDER CHRISTOPH
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<b>Experimental Team:</b> DREES Jan Yvo
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<b>Samples:</b>	SmFeO <sub>3</sub> Sr <sub>3</sub> YCo <sub>4</sub> O <sub>10.5</sub>
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Instrument	Req. Days	All. Days	From	To
D9	10	10	26/07/2012	30/07/2012
			26/11/2012	02/12/2012

<p><b>Abstract:</b></p> <p>Very recently room-temperature multiferroic properties of SmFeO<sub>3</sub> have been reported [6]. Density functional theory calculations explain the ferroelectricity in this exciting material with the calculated canted antiferromagnetic structure (T<sub>N</sub>~670 K). This renders single crystal neutron measurements highly desirable now. Despite the strong neutron absorption cross-section of the Sm ion we found conditions where the total absorption is comparably small. This could be realized by an optimized sample geometry and the use of 0.511 Å hot neutrons in a measurement at the D9 diffractometer. Our single crystal neutron measurements reveal an exciting first result: the spin configuration calculated in Ref. [6] can not be correct, since we observe quite strong magnetic intensities that should be absent for this calculated magnetic structure. We obtain a different magnetic symmetry and our magnetic structure is able to explain the ferroelectric properties with the inverse Dzyaloshinskii-Moriya interaction as well. After our first measurement we propose to perform a more detailed measurement at room-temperature now and also a 2nd measurement in the high-temperature magnetic phase (T<sub>SR</sub> = 480 K).</p>
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## Experimental report for D9 measurement of multiferroic SmFeO<sub>3</sub>

SmFeO<sub>3</sub> crystallizes in the well-known perovskite structure [1]. Very recently, the discovery of ferroelectric polarization in SmFeO<sub>3</sub> has been reported concomitantly with the onset of antiferromagnetic ordering at  $T_N \sim 670$  K in this compound [1]. The origin of its multiferroic properties above room-temperature i.e. the origin of the ferroelectric polarization in SmFeO<sub>3</sub> is highly debated [2–3]. A further spin-reorientation transition occurs at  $T_{SR} \sim 480$  K in this material [1]. In the initial publication an inverse Dzyaloshinskii-Moriya interaction based mechanism has been reported to be the driving force of the ferroelectric properties of SmFeO<sub>3</sub> [1]. The underlying magnetic structure was calculated by ab initio calculations [1]. However, it has been demonstrated that this calculated  $k=0$  magnetic structure with magnetic ions located at inversion centers can not be responsible for a spin-orbit-coupling driven ferroelectric polarization by  $S_i \times S_j$  in this material since inversion symmetry will not be broken [2]. Finally, a different alternative magnetoelastic mechanism based on  $J S_i \cdot S_j$  exchange-striction has been proposed to be responsible for the ferroelectric polarization in SmFeO<sub>3</sub> [3]. In order to elucidate the physical properties of this highly controversially discussed SmFeO<sub>3</sub>-system we studied this intriguing system experimentally now.

In this D9 measurement we were able to measure the magnetic structure of SmFeO<sub>3</sub> *for the first time* experimentally by neutron diffraction despite the highly neutron absorbing properties of the Sm ion ( $\sim 5900$  barn for 2200 m/s neutrons). *This was achieved by the high incident neutron energy available at the D9 diffractometer as well as by the optimized sample geometry chosen.* The magnetic structure of SmFeO<sub>3</sub> has been measured in the two different magnetic phases i.e. at 300 K below and at 515 K above the spin-reorientation transition  $T_{SR}$ .

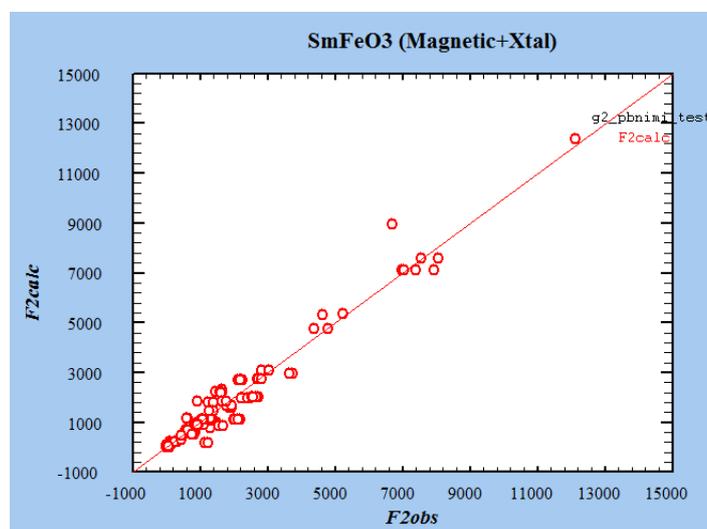
156 reflections have been measured at 300 K. The R-values of our magnetic and crystal structure refinement amount to  $R(F^2)$ : 18.7% and  $R_w(F^2)$ -factor : 21.1%. Other magnetic structures could be excluded. **Fig. 1** shows the  $F_{calc}/F_{obs}$  plot indicating the goodness of our fit for our “300 K refinement”. The finally obtained spin structure is shown in **Fig. 2**. This results already now prove that the interpretations in Ref. [1] can not be correct.

In a similar manner the 515 K spin structure could be obtained by the measurement of 156 reflections using a furnace (see **Fig. 3 & 4**). Also the “515 K spin structure” contradicts the whole interpretation presented in Ref. [1].

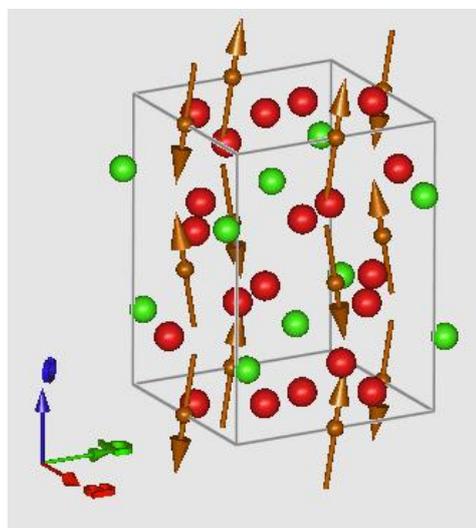
### REFERENCES:

- [1] J.-H. Lee et al., *Phys. Rev. Lett.* **107**, 117201 (2011)  
[2] R. D. Johnson et al., *Phys. Rev. Lett.* **108**, 219701 (2012)

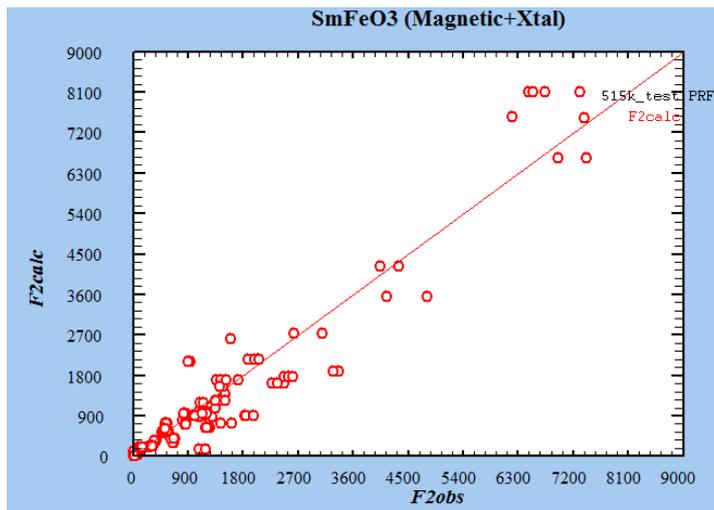
- [3] J.-H. Lee et al., *Phys. Rev. Lett.* **108**, 219702 (2012)



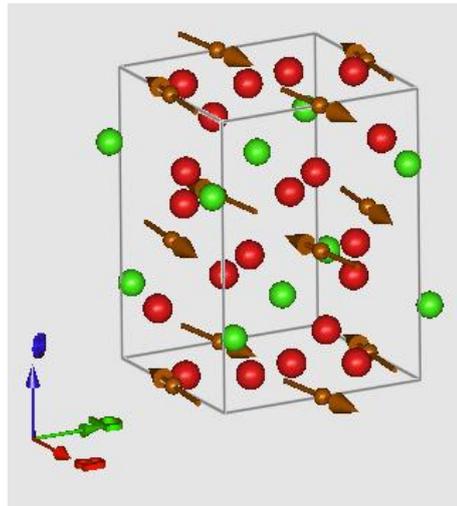
**Figure 1:** Resulting  $F_{calc}/F_{obs}$  plot of the magnetic structure refinement of SmFeO<sub>3</sub> at 300 K.



**Figure 2:** Magnetic structure of SmFeO<sub>3</sub> at 300 K. ( $Fe^{3+}$ -spins indicated by the arrows; red: O-ions; green: Sm-ions.)



**Figure 3:** Resulting  $F_{calc}/F_{obs}$  plot of the magnetic structure refinement of  $\text{SmFeO}_3$  at 515 K. Probably due to the use of a furnace the scattering is larger.



**Figure 4:** Magnetic structure of  $\text{SmFeO}_3$  at 515 K. ( $\text{Fe}^{3+}$ -spins indicated by the arrows; red: O-ions; green: Sm-ions.)