

# Experimental report

28/03/2017

**Proposal:** 5-41-877

**Council:** 4/2016

**Title:** Magnetic structure of  $\text{UnRhIn}_{3n+2}$  ( $n = 1, 2$ ) compounds

**Research area:** Physics

**This proposal is a continuation of 5-41-808**

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**Samples:**  $\text{U}_2\text{RhIn}_8$

$\text{URhIn}_5$

Instrument	Requested days	Allocated days	From	To
D10	8	4	16/11/2016	21/11/2016

## Abstract:

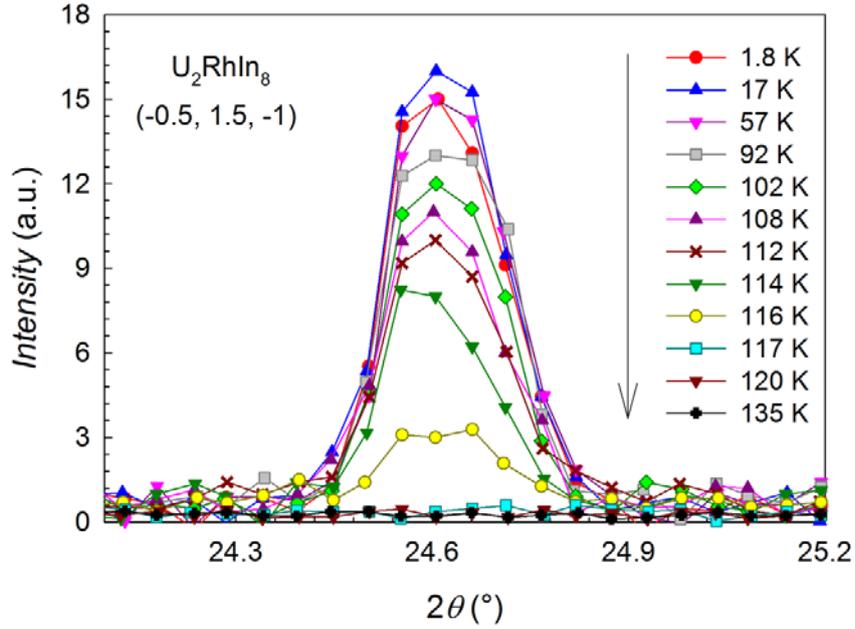
The recently discovered compound  $\text{U}_2\text{RhIn}_8$  belongs to the family with chemical composition  $\text{RnTmIn}_{3n+2m}$ , where R is a lanthanide or actinide and T a transition element. The materials grow with a tetragonal unit cell that can be viewed as n-layers of  $\text{RIn}_3$  stacked sequentially along the c-axis with m-layers of  $\text{TIn}_2$ . This layered structure allows investigating the influence of stacking on the magnetic and superconducting properties within a certain family. With  $\text{U}_2\text{RhIn}_8$  and  $\text{URhIn}_5$  we can study the effect of exchanging the localized 4f-electrons from the Ce-analogous with the more delocalized 5f-electrons. The cubic compound  $\text{UIn}_3$  orders antiferromagnetically at  $T_N = 88\text{K}$ .  $\text{U}_2\text{RhIn}_8$  and  $\text{URhIn}_5$  order at high temperatures,  $T_N = 117\text{K}$  and  $98\text{K}$ , respectively. The propagation vector of  $\text{UIn}_3$  equals  $(1/2, 1/2, 1/2)$ . The same wave vector is found for  $\text{URhIn}_5$ , while our recent experiment on D10 (proposal 5-41-808) revealed for the  $\text{U}_2\text{RhIn}_8$  compound a propagation vector of  $(1/2, 1/2, 0)$ . This proposal is a continuation of proposal 5-41-808. We like to determine the magnetic moment of the U in  $\text{U}_2\text{RhIn}_8$  as well as the temperature evolution of AFM state in  $\text{U}_2\text{RhIn}_8$  and  $\text{URhIn}_5$  to study the influence of the  $\text{RhIn}_2$ -layer.

# ILL Report: Magnetic structure of $U_nRhIn_{3n+2}$ ( $n = 1, 2$ ) compounds

**Scientific background:** Remarkably, many intriguing phenomena discovered in condensed matter physics during the past decades were found in intermetallic lanthanide ( $4f$ ) or actinide ( $5f$ ) compounds exhibiting tetragonal crystal structure. Examples are heavy fermion superconductivity in  $CeCu_2Si_2$  [1], hidden order in  $URu_2Si_2$  and quantum criticality beyond the Landau-Ginzburg-Wilson paradigm in for instance  $YbRh_2Si_2$ ,  $CeRhIn_5$  and  $Ce_2RhIn_8$  [2-4]. The lattice-type naturally inhabits a certain degree of frustration of the  $4f$  or  $5f$ -moments and the electronic dimensionality is often ascribed being less than 3-dimensional akin to layered structures. Both, frustration and dimensionality, play a crucial role in the description of earlier mentioned phenomena. In that sense, it is interesting to explore the effect of  $f$ -  $s,p,d$  electron hybridization on similar type of compounds which can be achieved by replacing the  $4f$ -electron element (Ce, Yb) by a  $5f$  element. (U). In the series,  $CeIn_3$ ,  $CeRhIn_5$  and  $Ce_2RhIn_8$  the Ce  $4f$ -electron is strongly localized. The compounds order antiferromagnetically with  $T_N$  of 10.2K, 3.8K and 2.8K respectively, and become superconducting under hydrostatic pressure.  $UIn_3$  has been exhaustively investigated. The compound orders antiferromagnetically (type-II AFM) at  $T_N = 88K$  with propagation vector  $(1/2, 1/2, 1/2)$  and moments aligned along  $\langle 110 \rangle$ -axis. Interestingly, NQR experiments reveal that the U-moments have localized character. Until recently solely  $U_nTGa_{3n+2}$  ( $T$ : transition element) compounds exist including only two so-called 2-1-8 compounds,  $U_2FeGa_8$  and  $U_2RhGa_8$  both being paramagnetic.  $URhIn_5$  was the first with In as  $p$ -element to be synthesized. In addition Bartha *et al.* [5] succeeded in preparing  $U_2RhIn_8$ , another  $5f$ - analogue in the Ce-series. Both compounds order AFM with Néel temperatures  $T_N = 117K$  (98K) for  $U_2RhIn_8$  ( $URhIn_5$ ). A recent NMR study on  $URhIn_5$  suggested that AFM here is driven by itinerant  $5f$  electrons and has propagation vector  $\mathbf{k} = (1/2, 1/2, 0)$  [6]. This proposed wave vector was confirmed by a recent neutron experiment on PANDA-FRM II [7]. The size of the magnetic moment yields  $\mu = 1.65 \mu_B/U$ . However, little is known about the magnetic properties of  $U_2RhIn_8$ . The current experiment focuses on the magnetic structure determination of  $U_2RhIn_8$ , to find out the magnetic moment on the Uranium and to follow the evolution of the AFM state with temperature in order to clarify the character of the magnetic ordering (2D or 3D). The experiment was a follow up experiment of ILL proposal 5-41-908 where already we were able to determine the propagation vector being  $(1/2, 1/2, 0)$ .

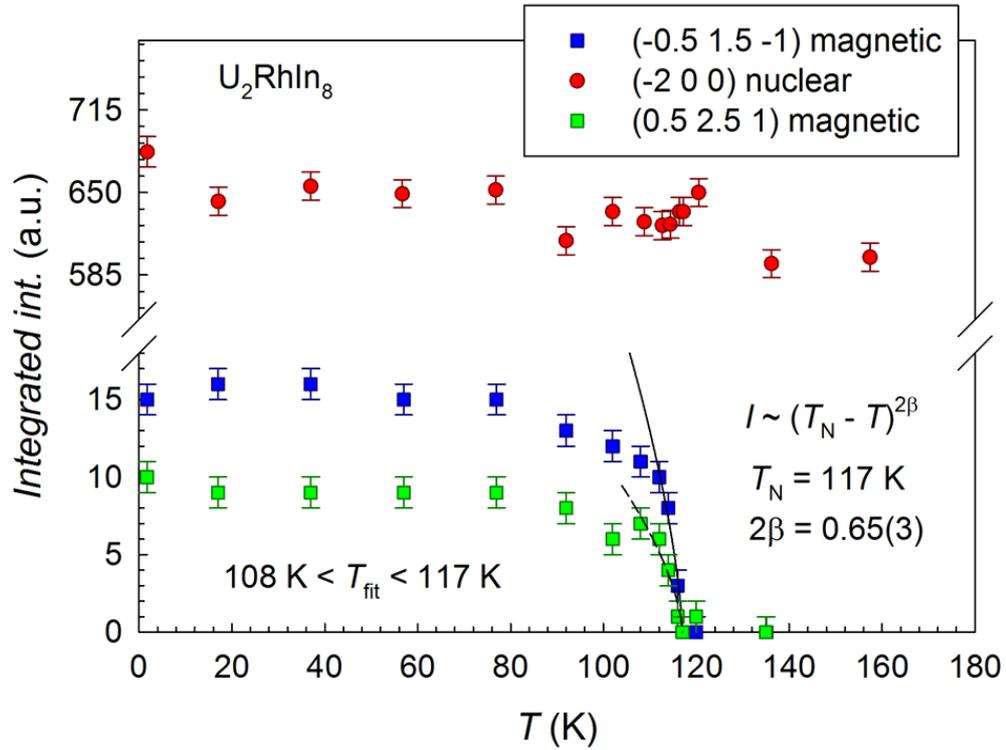
**Experimental:** Single crystals of  $U_2RhIn_8$  have been grown from In-flux. The samples are plate-like with  $c$ -axis perpendicular to the plate. Typical dimensions are  $1 \times 1 \times 0.5 \text{ mm}^3$  (mass  $\sim 1 \text{ mg}$ ). Prior to the neutron experiment, the samples were subjected to thermodynamic and transport experiments [5]. One sample was selected for further study on which we performed a Laue neutron single crystal diffraction experiment utilizing CYCLOPS instrument at ILL, Grenoble first. Scans were taken at  $T = 2 \text{ K}$  (AFM state) and  $130 \text{ K}$  (paramagnetic state). The presence of additional magnetic Bragg reflections in the  $2 \text{ K}$ -scan indicated that the ground state of  $U_2RhIn_8$  is AFM. The respective Laue patterns were further analyzed using Esmeralda software. By this we were able to confirm that the propagation vector in  $U_2RhIn_8$  indeed is  $\mathbf{k} = (1/2, 1/2, 0)$ . Subsequent neutron experiment was conducted on D10 in order to resolve the magnetic structure. The neutron wavelength equals  $\lambda = 2.36 \text{ \AA}$  and measurements were done down to  $T = 1.8 \text{ K}$ .

**Results:** We measured the intensity of two magnetic reflections  $(-0.5, 1.5, -1)$  and  $(0.5, 2.5, 1)$  at various temperatures between  $1.8 \text{ K}$  and  $135 \text{ K}$ . Figure 1 exemplarily presents the temperature dependence of the measured intensity  $(-0.5, 1.5, -1)$ . Clearly, intensity which relates to the ordered moment on the U-ions sets in below  $T_N$  ( $117 \text{ K}$ ) and increases upon lowering temperature. In order to follow the evolution of the ordered moment of the U-ion the integrated intensity has been calculated and plotted versus temperature in Fig. 2.

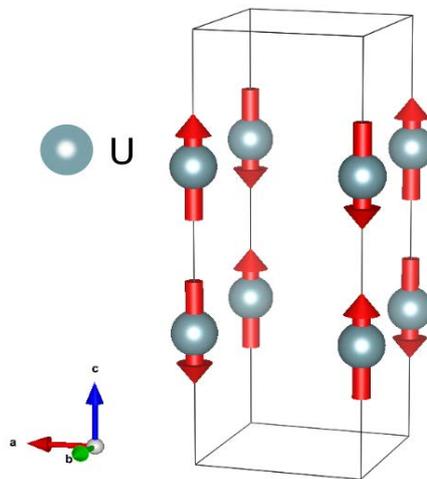


**Fig.1:** Temperature dependence of the intensity on the magnetic reflection  $(-0.5, 1.5, -1)$ . Error bars are within symbol size

Below approximately 80 K, the intensity shows no further increase. Figure 2 includes the nuclear reflection at reciprocal space position  $(-2, 0, 0)$ . The integrated intensity of the nuclear reflection is almost constant over the entire measured temperature range demonstrating the absence of a  $\mathbf{k}_0 = 0$  ferromagnetic component to the magnetic structure in accordance with results on CYCLOPS. Near the transition temperature, the integrated intensity of the magnetic diffraction peaks may be fit to a simple power law,  $I \propto (T_N - T)^{2\beta}$  with  $T_N$  the Néel temperature and  $\beta$  is the order parameter critical exponent. The fits were performed in a small interval of  $108 \text{ K} < T < 120 \text{ K}$  yielding  $T_N = 117 \text{ K}$  in agreement with macroscopic measurements [5]. For the critical exponent  $\beta$  we obtained for both datasets a value of  $\beta = 0.325$ . The value points to a 3D-Ising ordering of the U-moments [8]. In the following we employed MAXMAGN software using the refined structure parameters of  $\text{U}_2\text{RhIn}_8$  and propagation vector  $\mathbf{k}$  as input. We obtained 10 1D irreducible representations associated with the  $\mathbf{k}$ , but only 6 of them are part of the global reducible magnetic representation of the  $2g$  Wyckoff site occupied by U-ions. All 6 subgroups assume unit-cell doubling in  $a$ -direction. The magnetic space groups  $P_C4/mbm$  (127) and  $P_C4/nbm$  (125) have the magnetic moment aligned  $\mu$  along  $[001]$ . The  $C_a m m a$  (67) and  $C_a m m m$  (65) require alignment along  $[100]$  and the remaining two subgroups  $P_B m n a$  (53) and  $P_B m m a$  (51) have  $\mu$  parallel to  $[110]$ . To further identify the magnetic structure we measured the intensities of magnetic satellites described by  $\mathbf{k}$  at  $T = 1.8 \text{ K}$ . From the 174 magnetic reflections only 34 were independent. For refinement we used FullProf and we revealed best agreement with our data considering a commensurate magnetic structure with moments arranged along  $[001]$ , i.e., tetragonal  $c$ -axis. The agreement factors yield  $RF = 13.4\%$  and  $RF^2 = 12.1\%$ . The magnetic structure is depicted in Fig. 3. It corresponds to the magnetic space group  $P_C4/nbm$  – maximal magnetic subgroup of the paramagnetic space group  $P4/mmm$ .



**Fig.2:** Temperature dependence of the integrated intensity of the magnetic reflections  $(-0.5, 1.5, -1)$  and  $(0.5, 2.5, 1)$ . For comparison the integrated intensity of the nuclear reflection on  $(-2, 0, 0)$  is added. The solid (dashed) black line is power law fits in order to determine  $\beta$  from  $(-0.5, 1.5, -1)$  and  $(0.5, 2.5, 1)$ , respectively.



**Fig.3:** Magnetic structure of  $U_2RhIn_8$ . For clarity only the U-ions are displayed.

#### References:

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