

Experimental report

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Council: 10/2014

Title: Spin-lattice coupling in magnetoelastic 4d and 5d oxides

Research area: Physics

This proposal is a resubmission of 4-01-1384

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Samples: Ba₃BiIr₂O₉
Ba₃BiRu₂O₉
Ba₄BiIr₃O₁₂

Instrument	Requested days	Allocated days	From	To
IN4	7	7	19/06/2015	26/06/2015

Abstract:

We wish to investigate the giant magnetoelastic transition of Ba₃BiIr₂O₉, which shows an extremely rapid, large (1%) and anisotropic volume increase on cooling through a critical temperature $T^* = 72$ K, and isostructural Ba₃BiRu₂O₉. The volume increase is driven by a 5% increase in the Ir-Ir bond length at T^* , and is accompanied by a sharp drop in magnetic susceptibility. We believe the most likely explanation is the opening of a spin-gap, where the observed rapid decrease in susceptibility is related to the formation of local spin singlets (dimers). We previously collected TOF INS data for Ba₃BiIr₂O₉ and Ba₃BiRu₂O₉ on IN4, observing a spin-gap signal for the Ru compound at 33 meV. We now wish to explore this feature in greater detail and search for the gap in Ba₃BiIr₂O₉, which ab initio calculations indicate should have approximately half the energy. We also wish to search for evidence of a magnetoelastic wave associated with the phase transition.

Inelastic neutron scattering (INS) measurements were carried out on the direct-geometry, thermal-neutron time-of-flight spectrometer IN4C at the Institut Laue Langevin (Grenoble, France). An incident wavelength $\lambda_i = 1.11 \text{ \AA}$ ($E_i = 66.4 \text{ meV}$) was selected using a pyrolytic graphite monochromator. 20 g of polycrystalline $\text{Ba}_3\text{BiRu}_2\text{O}_9$ was sealed into a thin aluminium foil that was fixed to the cold tip of the sample stick of a standard orange cryostat. Measurements were performed at 100 K and 200 K, below and above $T^* = 176 \text{ K}$ respectively. The scattering function $S(Q, E)$ was measured in the neutron energy loss mode, in which the setting used in the down-scattering regime leads to a momentum transfer (Q) and energy transfer (E) extending up to 10 \AA^{-1} and 60 meV , respectively. Standard corrections including detector efficiency calibration and background subtraction were performed. The data analysis was done using ILL software tools.

Fig. 1 shows the color-coded Bose-factor corrected $S(Q, E)$ maps of the scattering intensity obtained on IN4C for $\text{Ba}_3\text{BiRu}_2\text{O}_9$ above, and at three temperatures below, $T^* = 176 \text{ K}$. A spin-gap excitation peak is clearly seen to emerge at low Q below T^* . Fig. 2 shows scans taken through the peak at $E_f = 14.87 \text{ meV}$ and $S_2 = 15.5^\circ 2\theta$ ($|Q| = 2.60 \text{ \AA}^{-1}$ at $E = 36 \text{ meV}$) at 3 and 200 K. The spin-gap peak was fitted to a Gaussian with respect to energy transfer to yield a spin-gap value $E_{exp} = 36 \pm 1 \text{ meV}$. Fig. 2(b) shows the temperature dependence of this peak intensity, fit to an order parameter $((T_c - T)/T_c)^{2\beta}$, yielding $T_c = 175 \text{ K}$ and $\beta = 0.204$, perfectly consistent with $T^* = 176 \text{ K}$. An analogous spin-gap excitation was not observed for $\text{Ba}_3\text{BiIr}_2\text{O}_9$; however, this is unsurprising given the very high neutron absorption cross-section of Ir and the extremely rapid fall-off of the Ir^{4+} magnetic form-factor, both of which obscure magnetic scattering intensity.

From previous experiments, no magnetic Bragg peaks were observed in neutron powder diffraction data for $\text{Ba}_3\text{BiRu}_2\text{O}_9$ [8] below T^* , indicating a lack of (observable) long range magnetic ordering. A model for diffuse scattering for liquids/gases etc., which is generally used for dimer systems, was therefore used to fit the INS data. Assuming the dimer state can be described by the wavefunctions $|s_1, s_2, M\rangle$ at the energy transfer corresponding to the spin gap, the intensity of scattered neutrons is given by [25]:

$$I \propto \exp\left(-\frac{\Delta_0}{k_B T}\right) \exp(-2W) f^2(Q) \mathbf{M} \left(1 - \frac{\sin(QR)}{QR}\right) \quad (1)$$

where Δ_0 is the spin gap, $\exp(-2W)$ is the Debye-Waller

factor, $f(Q)$ the magnetic form factor and \mathbf{M} is the matrix element describing the transition, and R is the metal-metal distance. In this case, we wish to study whether the effect is due to intra-dimer or inter-dimer alignment. The last term describes interference effect as a result of the metal-metal distance in the dimer.

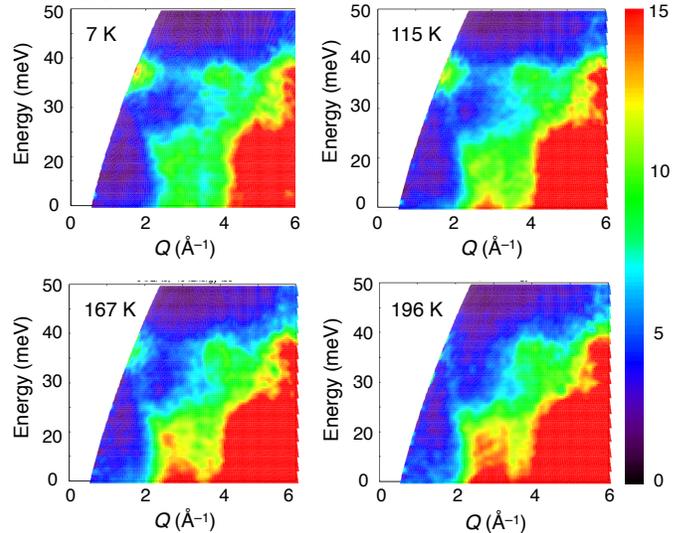


FIG. 1. Color-coded Bose-factor corrected $S(Q, E)$ maps of the inelastic neutron scattering of $\text{Ba}_3\text{BiRu}_2\text{O}_9$ taken on Merlin at ISIS above, and at three temperatures below, $T^* = 176 \text{ K}$. A magnetic excitation emerges below T^* at low- Q and an energy transfer of 36 meV . The color scale at right shows intensity (arbitrary units).

An integrated intensity slice was taken over the energy transfer range $32\text{--}42 \text{ meV}$ at 7 K and fitted to Equation (1) from $Q = 1.6$ to $Q = 2.4$. Larger Q values were excluded because the phonon background became significant. The form factor for Ru^{5+} [26] was used to perform the analysis, as there is no available experimental form factor for Ru^{4+} , and a phonon background was modeled of the form $xQ^2 + y$. Two alternative fits to determine the Ru-Ru distance were performed: one starting from the experimental intra-dimer Ru-Ru distance of 2.6 \AA ; and another starting from the (average) experimental inter-dimer distance of 5.9 \AA [8]. The former converged to an intra-dimer distance of $R = 2.61 \pm 0.01 \text{ \AA}$, and the latter to $R = 5.5 \pm 0.2 \text{ \AA}$. Note that the second discrete peak near $Q = 4$ is a phonon peak due to the aluminum sample holder. The intra-dimer model, shown in Fig. 3, produced the better fit to the low Q intensity. However, it should be acknowledged that without very low Q data, the inter-dimer model cannot be conclusively ruled out.

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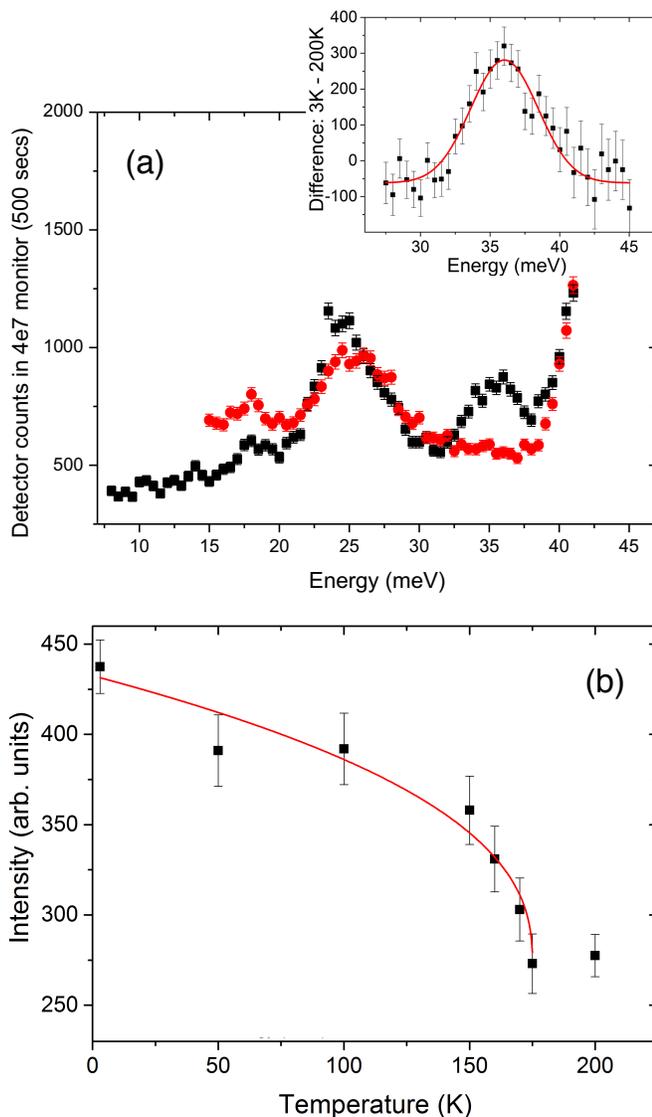


FIG. 2. Inelastic neutron scattering from $\text{Ba}_3\text{BiRu}_2\text{O}_9$, measured on Taipan at ANSTO. (a) Shows low- Q scans above (200 K, red) and below (3 K, black) $T^* = 176$ K. The inset shows the difference (3 K - 200 K) fit to a Gaussian centred at $E_{exp} = 36.03 \pm 0.14$ meV. (b) The intensity of this peak as a function of temperature, fit to an order parameter (see text for details).

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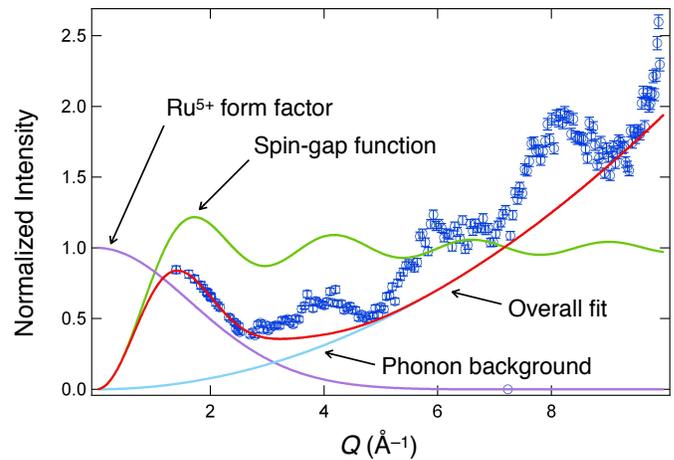


FIG. 3. Integrated intensity over the energy transfer range 32-42 meV of the Merlin data at 7 K, fit to equation (1). Components of the fit are labelled (see text for details). The oscillations at higher Q in the experimental data (blue circles) are non-magnetic (phonon) peaks due partly to the aluminum sample holder.

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